

Properties of Linear Algebra Applicable to Quantum Computing

Patrick Dreher CSC591 / ECE592 – Fall 2020

OUTLINE:

Mathematics and Physics Concepts Required to Describe a Quantum Computing System

- Introduce relevant mathematics applicable to quantum computing
- Introduce quantum mechanics axioms that describe observed physics
- Combine relevant mathematics with the physics of quantum mechanics to allow one to formulate
 - Design programming primitives that can aggregate into quantum computing algorithms
 - Program algorithms for implementation on quantum simulators and HW platforms
 - Configure and run these algorithms on quantum computing simulators and quantum computing hardware platforms

NC STATEUNIVERSITY []

Theory of Quantum Mechanics Describes Behavior Observed in a Non-classical World

- In order to properly design quantum computing algorithms programs and physical devices to performs calculations based on the observables axioms of quantum mechanics one must
 - Understand the quantum mechanical properties and measured observables
 - Use appropriate mathematics to properly describes these properties and observables
- Quantum theory is a mathematical model of the physical world at a scale where the size of the observation being made are of the same order of magnitude as the size of the object being observed
 - Measuring the speed of a car moving on a road (classical)
 - Measuring the bound state energy of an electron in a hydrogen atom (quantum mechanical)
- Many of the observed behaviors of the physical world at the quantum level have no analogs in people's everyday (classical) experiences

Building a Rigorous Mathematical Foundation for Describing Quantum Computing

Utilize the Mathematics of Linear Algebra to Represent Quantum Computing Processes

Mathematics Applicable to Quantum Computing

- The axioms of quantum mechanics are well described by the mathematics of linear algebra
- Linear algebra concepts for describing a quantum computing system
 - Any given system is identified with some finite- or infinite-dimensional Hilbert space
 - The system can be mathematically represented as states
 - Pure states correspond to vectors of norm 1
 - This norm 1 set of all pure states can be graphically represented/visualized by a unit sphere in a Hilbert space

- <u>Vector Space</u>: A vector space is a collection vectors, which may be added together and multiplied by scalar quantities and still be a part of that same collection of vectors
- The <u>adjoint</u> a[†] is the complex conjugate transpose of a column vector "a" and is sometimes called the Hermitian conjugate
- The space is complete as expressed by the norm

 $||a|| = (\langle a | a \rangle)^{1/2}$

Linear Dependence and Linear Independence

 A set of vectors is said to be linearly dependent if one of the vectors in the set can be defined as a linear combination of the others

 A set of vectors is said to be linearly independent if no vector in the set can be written according to the previous statement

Basis Vectors

A set of elements (vectors) in a vector space V is called a basis, or a set of basis vectors, if the vectors are

- linearly independent
- every vector in the vector space is a linear combination of this set

A basis is a linearly independent spanning set

Review Basic Linear Algebra Concepts Properties and Definitions of a Vector Space

- Given a vector space V containing vectors A, B, C the following properties apply
 - Commutativity [A+B=B+A]
 - Associativity of vector addition [(A+B)+C=A+(B+C)]
 - Additive identity [0+A=A+0=A] for all A
 - Existence of additive inverse: For any A, there exists a (-A) such that A+(-A)=0

Properties and Definitions of a Vector Space

- Given a vector space V containing vectors A, B, C the following properties apply
 - Scalar multiplication identity [1·A=A]
 - Given scalars r and s
 - Associativity of scalar multiplication [r(sA)=(rs)A]
 - Distributivity of scalar sums [(r+s)A=rA+sA]
 - Distributivity of vector sums [r(A+B)=rA+rB]

Hilbert Space

• <u>Hilbert Space</u>: A Hilbert space is a vector space G that has an inner product <u,v> such that a norm

$$|u| = \sqrt{\langle u, u \rangle}$$

turns G into a complete metric space. (A complete metric space is a metric space in which every Cauchy sequence is convergent)

- A Hilbert Space with vectors a and b may be defined over either the real or complex numbers with an inner product <b|a>
- A Hilbert Space maps an ordered pair of vectors to the complex numbers with the following properties
 - Positivity <a|a> is greater than 0 for |a> greater than 0
 - Linearity $<c|(\alpha|a> + \beta|b>) = \alpha <c|a> + \beta <c|b>$ where α and β are complex constants
 - Skew symmetry <b|a> = (<a|b>)*
- An example of a finite-dimensional Hilbert space is an n-dimensional vector space of the complex numbers Cⁿ with <u,v> as the vector dot product of u and v* (complex conjugate)

Dirac "bra" and "ket" Notation

Dirac "ket" notation |a> represents a column vector \vec{a}

$$a \ge = \begin{pmatrix} a_1 \\ a_2 \\ \vdots \\ a_n \end{pmatrix}$$

a Dirac "bra" notation <a

$$| < a | = (a_1^* \quad a_2^* \quad \dots \quad a_n^*)$$

The **transpose a^T** of a column vector a is a row vector

25-Aug 2020 27-Aug-2020

Examples of Normalized Vectors in Dirac Notation

$$\mathbf{a} = \frac{1}{\sqrt{2}} \left[|0\rangle + |1\rangle \right] = \frac{1}{\sqrt{2}} \left[\begin{pmatrix} \frac{1}{0} \end{pmatrix} + \begin{pmatrix} \frac{0}{1} \end{pmatrix} \right] = \left(\frac{\frac{1}{\sqrt{2}}}{\frac{1}{\sqrt{2}}} \right)$$

$$|b\rangle = \left[\frac{3}{5}|0\rangle - \frac{4}{5}|1\rangle\right] = \frac{3}{5}\left(\frac{1}{0}\right) - \frac{4}{5}\left(\frac{0}{1}\right) = \left(\frac{\frac{3}{5}}{\frac{-4}{5}}\right)$$
$$|c\rangle = \frac{3i}{5}|0\rangle - \frac{4i}{5}|1\rangle = \frac{3i}{5}\left(\frac{1}{0}\right) - \frac{4i}{5}\left(\frac{0}{1}\right) = \left(\frac{\frac{3i}{5}}{\frac{4i}{5}}\right)$$



Mathematical Representation of Binary States and Superposition

- A binary state (classical bit) defines a state by
- values of either "0" or "1" ("on" or "off")





NC STATEUNIVERSITY Y

15

Mathematical Representation of Bits, Qubits and Superposition

- A classical bit defines a state by values of either "0" or "1" ("on" or "off")
- A quantum bit (qubit) can also have a state of "0" or "1" but it can also have a possibility of being described by additional states



NC STATEUNIVERSITY []

Bits, Qubits and Superposition A classical bit defines a state by values

Mathematical Representation of

of either "0" or "1" ("on" or "off")

- A quantum bit (qubit) can also have a state of "0" or "1" and it can also have a possibility of being described by additional states
- Qubit can form a superposition state represented by a vector that is a superposition or linear combination of both a "0" or "1" |a> = α|0> + β|1> |α|² + |β|² = 1







A Qubit

• A qubit is a quantum system described by a two-dimensional Hilbert space, whose state can take any value of the form

 $|a\rangle = \alpha |0\rangle + \beta |1\rangle$

- We can perform a measurement that projects the qubit onto the basis {|0>, |1>} which will measure an outcome of |0> with probability $|\alpha|^2$ and an outcome of |1> with probability $|\beta|^2$ $|\alpha|^2 + |\beta|^2 = 1$
- Qubit can form a superposition state represented by a vector that is a superposition or linear combination of both a "0" or "1"
- The coefficients α and β also encode the relative phase that has physical significance

Basis Vectors for One Qubit

• In Dirac notation a qubit can be represented by

 $a = \alpha |0\rangle + \beta |1\rangle \qquad |\alpha|^2 + |\beta|^2 = 1 \text{ (modulus)}$

where α and β are complex coefficients

- α is the probability amplitude of measuring the |0> state and β is the probability amplitude of measuring the |1> state
- Common basis is $|0> = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $|1> = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$
- Probability to measure the |0> state is $|\alpha|^2$
- Probability to measure the |1> state is $|\beta|^2$
- Calculate <0|0> and <1|1> (gives an answer of a single number)

Constructing Matrices from Bras and Kets

If the bra and ket are placed in the opposite order

$$|0><0| = \begin{pmatrix}1\\0\end{pmatrix}(1 \quad 0) = \begin{pmatrix}1&0\\0&0\end{pmatrix}$$
$$|0><1| = \begin{pmatrix}1\\0\end{pmatrix}(0 \quad 1) = \begin{pmatrix}0&1\\0&0\end{pmatrix}$$
$$|1><0| = \begin{pmatrix}0\\1\end{pmatrix}(1 \quad 0) = \begin{pmatrix}0&0\\1&0\end{pmatrix}$$
$$|1><1| = \begin{pmatrix}0\\1\end{pmatrix}(0 \quad 1) = \begin{pmatrix}0&0\\0&1\end{pmatrix}$$

Outer products are a useful mechanism for writing matrices, especially unitaries because they capture state transformations

25-Aug 2020 27-Aug-2020

A Qubit

- Interpret a qubit from a geometric standpoint
- A $|0\rangle$ and $|1\rangle$ can represent opposite orientations of a vector in a three-dimensional space along a constructed set of axes in a coordinate system (example: z-axis with polar angle θ and azimuthal angle ϕ)
- A vector in this geometric construction can be transformed by rotations that express symmetries of the operations



Symmetries and Spatial Rotations

- Consider the case of a three dimensional rotation
- Given a vector ${\bf v}$ one can perform an infinitesimal rotation of ${\bf v}$ by d Θ about the axis specified by a unit vector
- Mathematically $R(\hat{n}, d\Theta) = I id\Theta \hat{n} \cdot \vec{J}$ where $\hat{n} = (n_1, n_2, n_3)$ are the components of angular rotations \vec{J}
- Basic properties of rotations can be written as commutation relations

$$[\mathbf{J}_{i}, \mathbf{J}_{j}] = \mathbf{J}_{i} \mathbf{J}_{j} - \mathbf{J}_{j} \mathbf{J}_{i}$$
$$[\mathbf{J}_{i}, \mathbf{J}_{j}] = i \varepsilon_{ijk} \mathbf{J}_{k}$$

where ε_{ijk} is the totally antisymmetric tensor and repeated indices are summed

Rotation Group

- Although the "defining" representation of the rotation group is three dimensional, the simplest nontrivial irreducible representation is two dimensional
- In this case there is a unique two-dimensional irreducible representation, up to a unitary change of basis.
- The generators of this rotation group are

$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \sigma_x$$
$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \sigma_y$$
$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \sigma_z$$

Properties of These Rotation Matrices

- Matrices are mutually anticommuting
- They mathematically can be squared to the identity

 $\sigma_i \sigma_j + \sigma_j \sigma_i = 2\delta_{ij}I \qquad (\hat{n} \cdot \vec{\sigma})^2 = n_i n_j \sigma_i \sigma_j = n_i n_j I$

• In spherical coordinates the pure density matrix can now be written

$$\rho(\hat{n}) = \frac{1}{2}(I + \hat{n} \cdot \vec{\sigma})$$
$$U(\hat{n}, \Theta) = \exp(i\frac{\Theta}{2}\hat{n} \cdot \vec{\sigma}) = I\cos(\frac{\Theta}{2}) - i\hat{n} \cdot \vec{\sigma}\sin(\frac{\Theta}{2})$$

Rotation Operators

• Write the exponential in terms of a series expansion

$$e^A = \sum_{k=0}^{\infty} \frac{1}{k!} A^k$$

• This gives rise to 3 useful classes of unitary matrices (rotation operators) when they are exponentiated

$$e^{-i\sigma \cdot n\frac{\theta}{2}} = \cos\left(\frac{\theta}{2}\right) - i\sin(\frac{\theta}{2})\sigma \cdot \hat{n}$$
$$R_{\hat{n}}(\theta) \equiv e^{-i\theta n \cdot \frac{\sigma}{2}} = \cos\left(\frac{\theta}{2}\right)I - i\sin\left(\frac{\theta}{2}\right)(n_x X + n_y Y + n_z Z)$$

Rotation Group

- The "defining" representation of the rotation group is three dimensional, but the simplest nontrivial irreducible representation is two dimensional
- In this case there is a unique two-dimensional irreducible representation, up to a unitary change of basis.
- The generators of this rotation group are

Pauli X – X –
$$\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \sigma_x$$

Pauli Y – Y – $\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \sigma_y$
Pauli Z – Z – $\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \sigma_z$
Pauli Spin Matrices have a special relationship in physics to particles that carry a property known as "spin"

Mathematical Representation of Many Different Basis States

- Represent combination of "0"s and "1"s in a way that many different values can be expressed
- Define $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ and $|1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$
- Can re-write $|a\rangle = \alpha |0\rangle + \beta |1\rangle$ as $|\alpha|^2 + |\beta|^2 = 1$

$$|a\rangle = e^{i\gamma} (\cos(\frac{\theta}{2})|0\rangle + e^{i\phi} \sin(\frac{\theta}{2})|1\rangle)$$

- This representation is visualized by states that lie of the surface of a Bloch sphere
- The Bloch sphere is a geometrical representation of the pure state space of a two-level quantum mechanical system

 $\hat{\mathbf{z}} = |0\rangle$

 $-\hat{\mathbf{z}} = |1\rangle$

https://en.wikipedia.org/wiki/Bloch sphere

Bloch Sphere

Figure from Wikipedia Bloch Sphere

 $\hat{\mathbf{x}}$

Matrices as Rotations Acting on Qubits

- Matrices describe the rotations that takes a qubit from an initial state to a transformed state
- These rotations that operate on a qubit are labelled as "gates"
- Because qubit states can be represented as points on a sphere, reversible one-qubit gates can be thought of as rotations of the Bloch sphere. (quantum gates are often called "rotations")
- Reversible one qubit gates viewed as rotations in this three dimensional representation

NC STATEUNIVERSITY

Construct Rotation Matrices From Bra and Ket Vectors

• The matrix representation of the expression $\sum_{i} |input_i| > < output_i|$

$$I = |0 > < 0| + |1 > < 1| = {1 \choose 0} (1 \quad 0) + {0 \choose 1} (0 \quad 1) = {1 \choose 0} (1 \quad 0)$$
$$X = |0 > < 1| + |1 > < 0| = {1 \choose 0} (0 \quad 1) + {0 \choose 1} (1 \quad 0) = {0 \choose 1} (1 \quad 0)$$
$$Z = |0 > < 0| - |1 > < 1| = {1 \choose 0} (1 \quad 0) - {0 \choose 1} (0 \quad 1) = {1 \choose 0} (1 \quad 0)$$
$$Y = iXZ = i {0 \choose 1} {1 \choose 0} {1 \choose 0} {1 \choose 0} = i {0 \choose 1} {-1 \choose 1} = {1 \choose 0} (1 \quad 0)$$
$$H = \frac{1}{\sqrt{2}} [(|0 > +|1 >) < 0| + (|0 > -|1 >) < 1|] = \frac{1}{\sqrt{2}} {1 \choose 1} {$$

Questions